

Lecture 1 - Macroscopic, entropy, and temperature - Summary

- **Main motivation of SM.**

An isolated system would exhibit *macroscopic properties* (like energy, magnetization etc.) that correspond to the highest number of *microscopic states*.

- **Entropy**

The natural log of the number of microscopic states, for a given set of Macroscopic measurables, $\{E\}$:

$$S(\{E\}) = \ln \Omega. \quad (1)$$

E could be total energy, total magnetization, volume, etc.

- **System of N electronic spins**

Magnetization:

$$M^z = \gamma \sum_i \hat{S}_i^z \quad (2)$$

$\gamma = g\mu_B = g\frac{e}{2m}$ is the gyromagnetic ratio, and \hat{S}_i^z is the z-direction spin of the i 'th electron.

M_z is a *conserved quantity*: it remains constant in an isolated system. Also, it is a prime example of a *macroscopic measurable*. The energy in a magnetic field H :

$$U = -HM_z = -H\gamma \sum_i \hat{S}_i^z. \quad (3)$$

Entropy with $N_\uparrow - N_{\text{downarrow}} = L$:

$$S(L) = \ln \Omega(L) \approx N \ln N - \frac{1}{2}(N+L) \ln \left(\frac{1}{2}(N+L) \right) - \frac{1}{2}(N-L) \ln \left(\frac{1}{2}(N-L) \right) + \mathcal{O}(\ln N) \quad (4)$$

As a function of U (for small U) it is:

$$\begin{aligned} S(U) &= N \ln N - \frac{1}{2}(N + \frac{2U}{H\gamma}) \ln \left(\frac{1}{2}(N + \frac{2U}{H\gamma}) \right) - \frac{1}{2}(N - \frac{2U}{H\gamma}) \ln \left(\frac{1}{2}(N - \frac{2U}{H\gamma}) \right) + \mathcal{O}(\ln N) \\ &\approx N \ln 2 - \frac{1}{2N} \left(\frac{2U}{H\gamma} \right)^2 \end{aligned} \quad (5)$$

This is called the **equation of state**.

The entropy was obtained using the Stirling formulae:

$$\ln n! = n \ln n - n + \mathcal{O}(\ln n). \quad (6)$$

- **Temperature**

$$\frac{1}{T} = \frac{\partial S(U)}{\partial U} \quad (7)$$

- **2nd law of Thermodynamics: Entropy always rises in non-equilibrium situations**

When two bodies come in contact and exchange energy:

$$\begin{aligned} \delta S_{total} &= \frac{\partial S_1(U_1)}{\partial U_1} \delta U_1 + \frac{\partial S_2(U_2)}{\partial U_2} \delta U_2 \\ &= \left(\frac{\partial S_1(U_1)}{\partial U_1} - \frac{\partial S_2(U_2)}{\partial U_2} \right) \delta U_1 \\ &= \left(\frac{1}{T_1} - \frac{1}{T_2} \right) \delta U_1 \geq 0 \end{aligned} \quad (8)$$

δU_1 has the same sign as $1/T_1 - 1/T_2$. So it is positive if $T_1 < T_2$. Heat flows from hot to cold.

- **First law of thermodynamics**

S is a **state function** of the macroscopic measurables of the systems. Generally, in addition to energy U , these will be volume V , number of particles N , and could anything else, which we mark X_i . As a state function, it is differentiable in all of its arguments. Therefore:

$$dS(U, V, N, \{X_i\}) = \frac{\partial S}{\partial U} dU + \frac{\partial S}{\partial V} dV + \frac{\partial S}{\partial N} dN + \sum_i \frac{\partial S}{\partial X_i} dX_i \quad (9)$$

Now:

$$\frac{\partial S}{\partial U} = \frac{1}{T}, \quad \frac{\partial S}{\partial V} = \frac{p}{V}, \quad \frac{\partial S}{\partial N} = -\frac{\mu}{T} \quad (10)$$

p is the pressure, and you see that if we change V , while keeping S constant, and also keeping N and the X 's constant, the energy changes by:

$$dU = -pdV. \quad (11)$$

The first law of thermodynamics:

$$TdS = dU + pdV - \mu dN + \dots \quad (12)$$

- **3rd law of Thermodynamics - Nernst Law: Entropy is a minimum at zero temperature**

The third law has a very simple quantum mechanical statement: The system is always at its quantum ground state at zero temperature. If the ground state is singly degenerate, the entropy at $T = 0$ is $S = 0$.

- **Negative temperature**

In the spin system discussed in lecture:

$$T = -\frac{N}{U} \left(\frac{H\gamma}{2} \right)^2 \quad (13)$$

Which is negative for $U > 0$. Although counterintuitive, negative temperatures were measured using NMR in LiF¹.

Some definitions:

- **Adiabatic process.** A process in which the entropy remains unchanged, and no absorption of heat occurs .
- **Extensive and intensive quantities.** If we magnify the system by factor b , and *extensive* quantity also grows by a factor b : $X \rightarrow bX$. Examples are total energy U , entropy S , volume V , number of particles N . *Intensive* quantities do not change under magnification. Examples: temperature T , pressure p , density $n = N/V$.

Typically for extensive quantities we use upper-case letters, and for intensive quantities lower-case letters (with the exception of temperature).

¹ E. M. Purcell, R. V. Pound, Phys. Rev. **81**, 279 (1951).